

**Time-dependent wave selection for information processing in excitable media**

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We demonstrate an improved technique for implementing logic circuits in light-sensitive chemical excitable media. The technique makes use of the constant-speed propagation of waves along defined channels in an excitable medium based on the Belousov-Zhabotinsky reaction, along with the mutual annihilation of colliding waves. What distinguishes this work from previous work in this area is that regions where channels meet at a junction can periodically alternate between permitting the propagation of waves and blocking them. These valvelike areas are used to select waves based on the length of time that it takes waves to propagate from one valve to another. In an experimental implementation, the channels that make up the circuit layout are projected by a digital projector connected to a computer. Excitable channels are projected as dark areas and unexcitable regions as light areas. Valves alternate between dark and light: Every valve has the same period and phase, with a 50% duty cycle. This scheme can be used to make logic gates based on combinations of OR and AND-NOT operations, with few geometrical constraints. Because there are few geometrical constraints, compact circuits can be implemented. Experimental results from an implementation of a four-bit input, two-bit output integer square root circuit are given.

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**I. INTRODUCTION**

The use of discrete propagating waves for information processing with Belousov-Zhabotinsky (BZ) excitable media began with the capillary tube logic gates constructed by Tóth and Showalter [1]. Since then many different schemes for information processing in BZ excitable media have been devised [2]. Steinbock *et al.* implemented a logic scheme in which a catalyst for the BZ reaction is printed onto regions of a planar medium [3]; waves are able to propagate in these regions, but not elsewhere, and the operation performed by a region is determined by the geometry of the region. Agladze *et al.* [4] carried out computer simulations and experimental demonstrations of a diode geometrical arrangement that permits a wave to travel in one direction, but not in the opposite direction. Other diode structures were reported in Refs. [5,6]. Motoike and Yoshikawa [7] simulated several logic circuit elements including AND, OR, and NOT gates, a time difference detector, and a memory element. An experimental implementation of a similar memory element was reported in Ref. [8]. Sielewiesiuk and Gorecki simulated a cross junction capable of acting as a coincidence detector and a switch [9] and also showed that when a train of pulses passes from one excitable region to another over a passive barrier, pulses will sometimes drop out of the train. The rate of dropout is determined by many different parameters [10].

The use of a light-sensitive ruthenium catalyst greatly simplifies the process of applying a pattern to a planar BZ excitable medium. The entire medium can be impregnated with the catalyst, but the reaction can be suppressed in regions where it is not wanted by illuminating those regions [11]. Typically a slide projector or a digital projector is used for

this [12–14]. Górecki [13] designed unary and binary counting circuits that are able to count the number of waves entering a channel. All of the components of the counters were tested in experiments, although the complete counters were not. Ichino *et al.* [14] showed that the behavior of a diode structure and a time difference discriminator can be altered by altering the level of illumination.

Whereas the construction of basic logical elements in BZ excitable media depends upon finding a mapping between logical operations and the behavior of waves in a medium, several works have demonstrated that some more complex information processing tasks can be carried out by finding a good mapping or analogy between the task and the dynamics of the medium. Kuhnert *et al.* projected an image directly onto a light-sensitive medium to perform an image processing task [12]. Steinbock *et al.* used circularly spreading waves to help find the shortest path through labyrinths [15], where the interior of the labyrinth was excitable, but the labyrinth walls and exterior of the labyrinth were not. Gorecka and Gorecki [16] simulated the operation of a bandpass filter, capable of outputting a signal whenever the time interval separating the arrival of two input pulses matched a constant value determined by the dimensions of the filter. Motoike and Yoshikawa [17] simulated a wave train comparator and subtractor that operate directly on trains of waves. Nagahara *et al.* [18] demonstrated a direction detector, capable of detecting the direction of origin of a spreading circular wave. Gorecki *et al.* [19] showed how an array of excitable paths can be used to make a distance detector by exploiting the phenomenon described in Ref. [10]. Gorecka and Gorecki [20] simulated a device analogous to an integrate-and-fire neuron and showed that it can effect a four-input Boolean threshold function much more simply than can be done using AND and OR gates.

Adamatzky explored the potential of BZ excitable media as a substrate for implementing architectureless

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collision-based computing schemes [21]. In this paradigm, exemplified in a cellular automaton model by Margolus in Ref. [22], the substrate does not contain any information about the computation to be carried out: It contains no structure that determines a specific function. This information is contained in the pattern and timing of events that are applied to the substrate. For certain levels of illumination, light-sensitive BZ excitable media have a subexcitable mode in which waves of a critical size travel forward and may either expand or contract under the influence of small perturbations [23]. Waves below this size slowly contract as they move forward. Waves above this size slowly expand. Collisions between subexcitable waves can be interpreted as logical operations [24,25]. Subexcitability combined with structured media was used in Ref. [26]. In Refs. [27,28], interactions between subexcitable waves were contained within circular disks to help control their behavior. Spiral waves were used in Ref. [29] as the basis for a glider-gun-like mechanism, capable of producing a continuous stream of discrete waves.

The variety of phenomena that are available in BZ excitable media and have been used for constructing information processing structures justifies the use of these media for exploring new ways of making information processing mechanisms. By investigating all of the possible ways in which these phenomena may be used for information processing, we equip ourselves with a tool set that can be applied when we encounter new media that share some of these phenomena.

The work in this paper has several things in common with the work of Steinbock *et al.* [3]. It makes use of the same two phenomena: The constant speed propagation of waves in BZ excitable media and the mutual annihilation of two wave fronts that meet. Differences in wave propagation time between paths from the inputs to the outputs of a region of excitable medium are exploited for implementing logical operations. Where this work differs from that of Steinbock *et al.* is that the region through which waves propagate is not a fixed pattern, but contains areas that alternate between being excitable and unexcitable in a simple periodic way. This is possible because the region is defined by a pattern projected using a digital projector. The ability to periodically switch areas from excitable to unexcitable and back is used to overcome one of the limitations of the work of Steinbock *et al.*, described below.

Steinbock *et al.* use the presence or absence of a wave in an input channel to represent Boolean 1 and 0 values at the inputs to a gate. All inputs to a gate must be applied at the same time. At the output of a gate, the synchrony or asynchrony of the arrival of waves in a pair of output channels is used to represent Boolean 1 and 0, respectively: If the waves arrive at both output channels at the same time, this represents 1; if the waves arrive at different times, this represents 0. One of the limitations of the gates of Steinbock *et al.* is that the timing of the output from a gate is dependent on the values of the inputs. Consequently, when gates are cascaded there is no way to be sure that all inputs at a gate downstream in the cascade arrive at the same time. This is illustrated in Fig. 1, which shows the  $z$  OR ( $x$  AND NOT  $y$ ) gate from Ref. [3]. Successive times for the case  $x = 0, y = 0,$  and  $z = 1$  are shown in Figs. 1(a) and 1(b). For this case, a 1 output will emerge from the gate after the time that it takes a wave to propagate from  $z$  to the output. However,

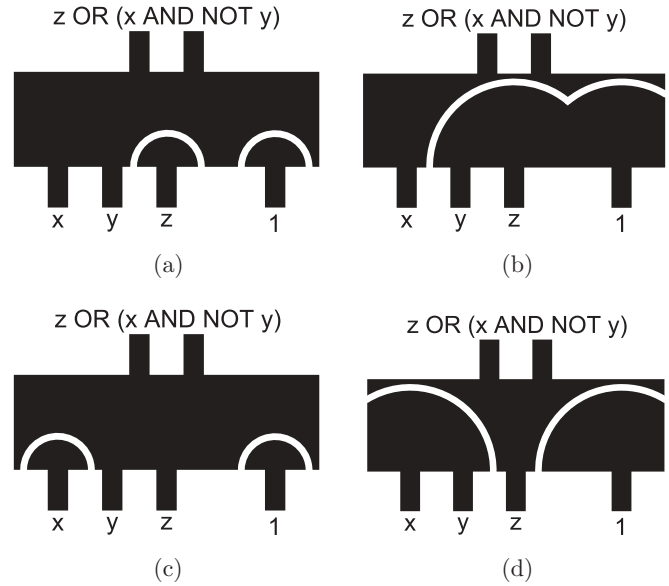


FIG. 1. The  $z$  OR ( $x$  AND NOT  $y$ ) gate from Ref. [3]. For the case  $x = 0, y = 0,$  and  $z = 1$  (a) shows waves that originated at the  $z$  input and the constant 1 input and (b) shows these waves at a later time. For the case  $x = 1, y = 0,$  and  $z = 0$  (c) shows waves that originated at the  $x$  input and the constant 1 input and (d) shows these waves at a later time. For both cases, waves arriving at the output channels are synchronized, but a comparison of (b) and (d) illustrates that for the latter case it takes longer for waves to reach the output than in the former case.

for the case  $x = 1, y = 0,$  and  $z = 0$ , shown in Figs. 1(c) and 1(d), a 1 output will emerge from the gate after the time that it takes a wave to propagate from  $x$  to the output, which is longer than in the previous case. Circuits implemented using the scheme described in this paper do not have this limitation. Outputs emerge from a circuit at a constant time after the inputs are applied.

## II. TIME-DEPENDENT WAVE SELECTION

### A. Fork structures

The path that led to the work in this paper began by observing the similarity between the behavior of toppling dominoes and the behavior of waves in channels of BZ excitable media. Domino fork structures can be used as the basis for making Boolean logic circuits [30,31]. A domino fork structure, shown in Fig. 2, has three lines of dominoes meeting at a junction. We call two of these lines arms and the third we call the stem. The end of each line is labeled with two

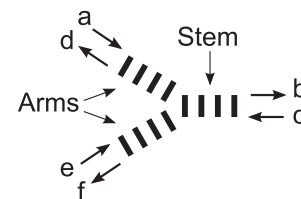


FIG. 2. Domino fork. Letters denote events. Input events ( $a, c,$  and  $e$ ) correspond to a domino toppling into the fork and output events ( $b, d,$  and  $f$ ) correspond to a domino toppling out of the fork.

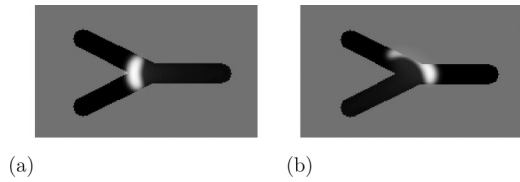


FIG. 3. Simulation of a fork structure in BZ excitable media. (a) Splitting at a fork junction in a simulation of subexcitable BZ-reaction media. (b) Directional propagation at a fork junction in a simulation of subexcitable BZ-reaction media.

different events: One event corresponds to a domino toppling into the structure and the other event corresponds to a domino toppling out of the structure. The behavior of the domino fork can be described in terms of the relationships between these events. Informally, we can say that either event *a* or event *e* (or both) will cause event *b* (unless event *c* occurs in the meantime) and that event *c* will cause events *d* and *f* (unless *a* or *e* occur in the meantime).

An important feature of fork behavior is that an event entering the fork from one arm will not influence the other arm; it will only cause an event to emerge from the stem of the fork. In the domino fork structure this behavior is obtained through directionally dependent propagation of a toppling domino wave: A domino can only influence other dominoes if they stand in the direction in which it topples. A fork structure with behavior based on directionally dependent wave propagation can also be implemented in BZ excitable media. A structure similar to this can be found in Refs. [7,26] and Fig. 3 shows the result of simulating a fork structure that depends upon subexcitability (as defined in Sec. I) for its operation. Subexcitability is used to prevent a wave that is propagating along one arm from entering the other arm. The obvious difference between a domino fork and a fork implemented in BZ excitable media is that whereas the domino fork cannot be used again once dominoes have toppled, the excitable medium fork can be used again and again.

Simulations of excitable media circuits were able to reproduce some of the behavior described in Ref. [31]. The success of these simulations led to attempts to implement them experimentally (using the same experimental setup described later in Sec. III B). Figure 4 shows snapshots from

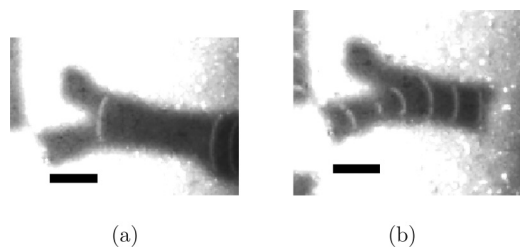


FIG. 4. Experimental implementation of a fork structure in BZ excitable media. The scale bar is 3 mm. The level of illumination in dark regions was estimated at  $0.5 \text{ mW cm}^{-2}$ . The level of illumination in light regions was estimated at  $6.5 \text{ mW cm}^{-2}$ . (a) Splitting at a fork junction in subexcitable BZ-reaction media [compare with Figure 3(a)]. (b) Directional propagation at a fork junction in subexcitable BZ-reaction media [compare with Fig. 3(b)].

an experiment exhibiting the same behavior as shown in Fig. 3. A video from this experiment can be found in Ref. [32].

Two problems were encountered in the course of conducting these experiments. The first was that the range of illumination levels in which the subexcitable regime can be obtained is both very narrow and also sensitive to experimental conditions. If the illumination level is too high, the medium is unexcitable and waves die off after a short distance. If the illumination level is too low, the medium is fully excitable and directionally dependent propagation cannot be obtained. The subexcitable illumination range was found to change from one location in the medium to another and also to change over time during an experimental session. It was therefore difficult to obtain a subexcitable regime across a large region for a long period of time.

The second problem was that harnessing subexcitability to obtain fork behavior places constraints on the geometry of a circuit made from forks: The arms of the fork must meet the stem within a certain range of angles and must extend for a certain distance away from the fork junction before they can turn a corner. These constraints limit the complexity of circuits that can fit into the area available for experiments (approximately  $5 \times 5 \text{ cm}^2$ ).

**B. Selecting waves based on their propagation time**

These problems led to a search for an alternative way to obtain the same behavior. Rather than use subexcitability to prevent a wave that is traveling along one arm from entering the other arm, we instead arrange the fork so that the time that it takes for a wave to propagate from arm to stem (or from stem to arm) is twice as long as the time that it takes for a wave to propagate from arm to arm. We then make sure that only waves taking the longer time are permitted to exit the fork.

Consider the T-shaped region of excitable medium shown in Fig. 5. This region is fully excitable: A wave entering the region will propagate throughout the region. A unit of length is defined in Fig. 5 and we define a unit of time to be the time that it takes a wave to propagate one unit of length. There are three terminals at the boundary of this region at which waves can be sent into the region or observed emerging from the region. We call the passage of a wave in a particular direction through one of these terminals an event and we label events using letters.

Table I shows which events emerge from the region at times  $t_0 + 1$  and  $t_0 + 2$  for all possible combinations of input events

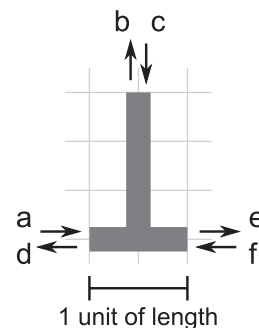


FIG. 5. A T-shaped region of excitable medium. The region is fully excitable, so a wave entering one terminal will propagate throughout the region. A unit of length is shown. A unit of time is defined as the time that it takes a wave to propagate one unit of length.

TABLE I. Behavior of the region in Fig. 5 for all possible combinations of input events  $a$ ,  $c$ , and  $f$  at time  $t_0$ .

Events at $t_0$	Events at $t_0 + 1$	Events at $t_0 + 2$
none	none	none
$a$	$e$	$b$
$c$	none	$d, e$
$f$	$d$	$b$
$a, c$	$e$	none
$a, f$	none	$b$
$c, f$	$d$	none
$a, c, f$	none	none

$a$ ,  $c$ , and  $f$  at time  $t_0$ . If we let an event at a specified time represent a Boolean 1 value and the absence of an event at a specified time represent a Boolean 0 value, then we can see from this table that if we ignore or somehow block events that occur at  $t_0 + 1$  and only pay attention to events that occur at  $t_0$  and  $t_0 + 2$  then this region can be used to implement two different Boolean operations.

(i) If  $c$  AND NOT ( $a$  OR  $f$ ) occur at  $t_0$  then both  $e$  and  $d$  will occur at  $t_0 + 2$ .

(ii) If ( $a$  OR  $f$ ) AND NOT  $c$  occur at  $t_0$  then  $b$  will occur at  $t_0 + 2$ .

The main insight that led to the results in this paper is that events that occur at  $t_0 + 1$  can be blocked by placing valves at each terminal and ensuring that the valves permit the passage of waves at  $t_0$  and  $t_0 + 2$ , but do not permit the passage of waves at  $t_0 + 1$ . In Sec. III B we will see that in an experimental setup where the circuit layout is projected by a digital projector onto the surface of a light-sensitive excitable medium, valves can be implemented by having valve regions of the projected circuit switch between dark and light periodically: When the valve regions are dark they permit the passage of waves; when they are light they inhibit the propagation of waves. We call this scheme time-dependent wave selection.

Figure 6 illustrates this for a T-shaped region in which two terminals have been designated as inputs  $a$  and  $b$  and the third terminal has been designated as an output. The output terminal outputs the function  $a$  OR  $b$  two time units after the inputs are applied. In this figure and those following a location labeled using the notation “ $a$  at  $t_0$ ” means that a Boolean 1 value at that location is represented by the occurrence of event  $a$  at time  $t_0$  and that a Boolean 0 value at that location is represented by the absence of event  $a$  at  $t_0$ .

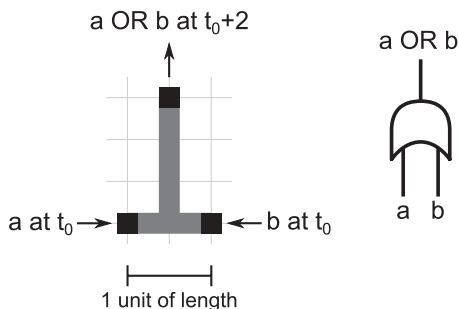


FIG. 6. A T-shaped region of excitable medium with valves at its terminals. This region implements a two-input OR gate.

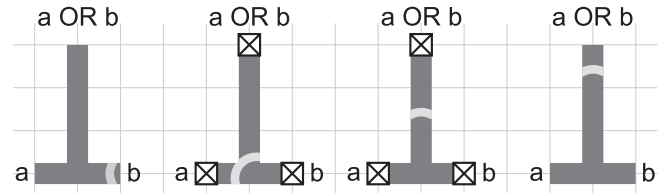


FIG. 7. A wave propagating from  $b$  to the  $a$  OR  $b$  output. Squares with crosses inside represent closed valves. The left-hand case shows the wave shortly after  $t_0$ . Subsequent cases show the progress of the wave every 0.5 time units until shortly after  $t_0 + 1.5$ . The wave never reaches  $a$  because it is blocked by a closed valve.

In order to permit some flexibility in the placement of valves and in order to tolerate minor deviations from perfect timing in simulations and experiments, valves alternate between being open (i.e., they permit the passage of waves) for a whole unit of time and then closed for a whole unit of time. Valves are open at  $t_0 + 2n - 0.5 < t \leq t_0 + 2n + 0.5$  for any integer  $n$  and closed at other times.

Figure 7 illustrates what happens to a wave that enters  $b$  at time  $t_0$ . The wave propagates throughout the excitable region, but it is unable to reach  $a$  because valves are closed around the time that the wave reaches  $a$ . However, by the time the wave reaches the  $a$  OR  $b$  output, valves are open and the wave is not blocked.

Fork behavior obtained through time-dependent wave selection overcomes the two problems that were encountered with trying to obtain the same behavior using subexcitability. Time-dependent wave selection uses the fully excitable regime and so it is not sensitive to experimental conditions. Additionally, it does not require any particular geometry: T shapes can be deformed in any way that preserves the wave propagation time between any pair of terminals.

Because time-dependent wave selection does not place many constraints on the geometry of a structure, it is easy to make structures having more than three terminals, so the behavior exhibited by the fork can be extended. The structure shown in Fig. 8 extends the OR-gate function from Fig. 6 by introducing an additional input, labeled  $c$ . Because this input is a unit distance away from all other terminals, a wave entering  $c$  at  $t_0$  will not emerge from any other terminal. However, it will annihilate any wave that enters at  $a$  or  $b$  at  $t_0$  and thus prevent waves from  $a$  or  $b$  from reaching the output. This structure implements the function  $(a$  OR  $b)$  AND NOT  $c$ .

The structure shown in Fig. 9 has five terminals: two inputs and three outputs. This structure implements the function  $a$

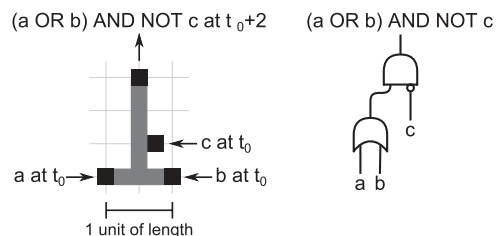


FIG. 8. A T-shaped region of excitable medium with an additional input. This structure implements the function  $(a$  OR  $b)$  AND NOT  $c$ .



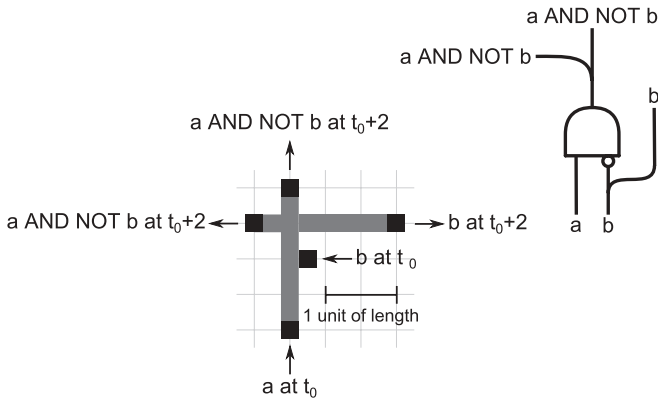


FIG. 9. Region of excitable medium that implements the function  $a$  AND NOT  $b$ . The  $a$  AND NOT  $b$  output emerges from two different terminals and the  $b$  input is also made available at an output.

AND NOT  $b$ , where the output emerges in two different locations and the  $b$  input is also made available at an output.

**III. FOUR-BIT SQUARE ROOT CIRCUIT**

The left-hand side of Fig. 10 shows a schematic diagram for a logic circuit that implements the four-bit integer square root function  $y = \lfloor \sqrt{x} \rfloor$ . In Fig. 10 and in later figures and tables,  $x_n$  is the  $n$ th binary digit of  $x$  and  $y_n$  is the  $n$ th binary digit of  $y$ . The truth table for the circuit is shown in Table II.

The right-hand side of Fig. 10 shows how the structures in Figs. 6, 8, and 9 can be used to implement this circuit. The logic gate symbols and excitable media channels in Fig. 10 are shaded to show the correspondence between the two and the key refers to the preceding figures in which the excitable media structures that are used here were introduced.

**A. Simulation**

A two-variable Oregonator equation [33] adapted for light sensitivity and with a diffusion term was used as an approximate model for the Belousov-Zhabotinsky reaction with applied illumination on a planar surface in which reactants

TABLE II. Truth table for a four-bit integer square root circuit.

$x$	$x_3$	$x_2$	$x_1$	$x_0$	$y = \lfloor \sqrt{x} \rfloor$	$y_1$	$y_0$
0	0	0	0	0	0	0	0
1	0	0	0	1	1	0	1
2	0	0	1	0	1	0	1
3	0	0	1	1	1	0	1
4	0	1	0	0	2	1	0
5	0	1	0	1	2	1	0
6	0	1	1	0	2	1	0
7	0	1	1	1	2	1	0
8	1	0	0	0	2	1	0
9	1	0	0	1	3	1	1
10	1	0	1	0	3	1	1
11	1	0	1	1	3	1	1
12	1	1	0	0	3	1	1
13	1	1	0	1	3	1	1
14	1	1	1	0	3	1	1
15	1	1	1	1	3	1	1

and products diffuse [34]:

$$\frac{\partial u}{\partial t} = \frac{1}{\epsilon} \left( u - u^2 - [fv + \phi(r,t)] \frac{u - q}{u + q} \right) + D_u \nabla^2 u,$$

$$\frac{\partial v}{\partial t} = u - v.$$

The variables  $u$  and  $v$  represent the local concentrations of activator and inhibitor, respectively. The parameter  $\epsilon$  sets up a ratio of time scale for the variables  $u$  and  $v$ ,  $q$  is a scaling parameter dependent on the rates of activation or propagation and inhibition, and  $f$  is a stoichiometric factor. The term  $\phi(r,t)$  is the rate of inhibitor production, which is proportional to the level of illumination. This varies over space  $r$  and for valve locations  $\phi(r,t)$  also varies with time  $t$ .

The system was integrated using the Euler method with a five-node Laplace operator for the diffusion term. The time step  $\Delta t$  was 0.001 and the grid point spacing  $\Delta x$  was 0.1 (which corresponds to 0.1 mm in the experimental setup). The following values were chosen so as to give behavior qualitatively similar to that observed in experiments:  $\epsilon =$

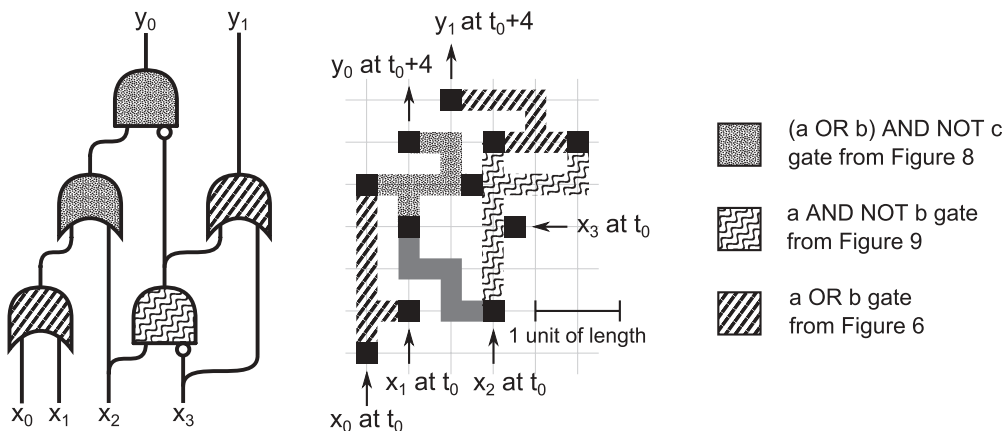


FIG. 10. Four-bit integer square root circuit. Here  $x_n$  is the  $n$ th binary digit of  $x$  and  $y_n$  is the  $n$ th binary digit of  $y$ . The logic gate symbols and excitable media channels are shaded to show the correspondence between the two. The key refers to the previous figures in which the excitable media structures that are used here were introduced.

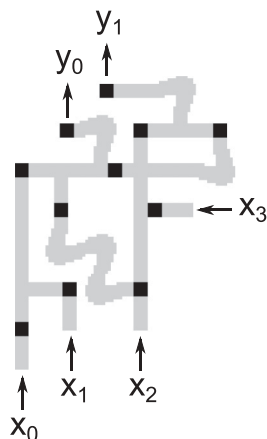


FIG. 11. Pattern for the four-bit square root circuit. Gray areas are areas of constant excitability. Dark areas are valves that alternate between excitable and unexcitable. White areas are unexcitable.

0.033,  $f = 1.4$ ,  $q = 0.002$ , and  $D_u = 0.067$ . For locations and times at which  $\phi(r,t)$  is excitable,  $\phi(r,t) = 0.055$ . For other locations and times,  $\phi(r,t) = 0.1$ . This system provides a good enough simulation of the experimental setup to help establish what types of behavior are likely to be achievable experimentally.

The pattern for the four-bit square root circuit is shown in Fig. 11. Each channel in Fig. 11 is 15 grid points wide. For gray areas of the pattern (excitable channels)  $\phi(r,t) = 0.055$ , for white areas (unexcitable regions)  $\phi(r,t) = 0.1$ , and for black areas (valves)  $\phi(r,t)$  is a square wave with period  $P$  and a 50% duty cycle:  $\phi(r,t) = 0.055$  when  $nP < t \leq nP + P/2$  for some integer  $n$ ; otherwise  $\phi(r,t) = 0.1$ . By measuring the number of simulation steps it took for a wave to travel from the valve nearest  $x_0$  to the next valve directly in line with  $x_0$ ,  $P$  was set to 15.5 (i.e., 15 500 simulation steps). Corners in the pattern mask are shaped so as to make the time for wave propagation along two channel segments at right angles to each other the same as when the two channel segments are placed together in a straight line.

Figure 13 at the end of Sec. III B shows snapshots from the simulation alongside experimental results for the cases  $x = 3, 6, \text{ and } 15$ . A video showing a simulation of all 15 cases can be found in Ref. [32].

## B. Experimental implementation

### 1. Experimental setup

Sodium bromate, sodium bromide, malonic acid, sulphuric acid, and tris(bipyridyl) ruthenium(II) chloride were purchased [Sigma-Aldrich, United Kingdom (UK)] and used as received unless stated otherwise. The catalyst free reaction mixture was freshly prepared in a 30 mL continuously fed stirred tank reactor (CSTR), which involved the *in situ* synthesis of stoichiometric bromomalonic acid from malonic acid and bromine generated from the partial reduction of sodium bromate. This CSTR in turn continuously fed a thermostatted open reactor with fresh catalyst-free BZ-reaction mixture in order to maintain a nonequilibrium state. The final composition of the catalyst-free reaction solution in the reactor was 0.42M sodium bromate, 0.19M malonic acid, 0.64M sulphuric

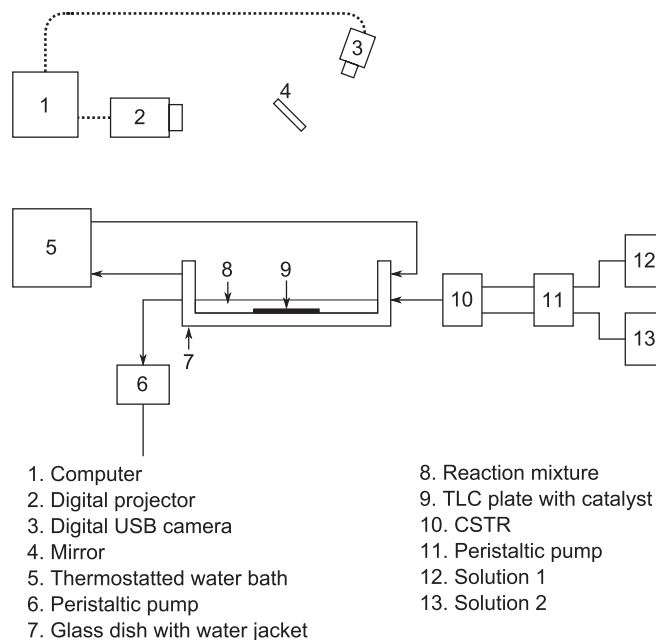


FIG. 12. Block diagram of the experimental setup.

acid, and 0.11M sodium bromide. The residence time was 30 min. The open reactor was surrounded by a water jacket thermostatted at 25 °C.

Peristaltic pumps (Model Sci-Q 400, Watson Marlow Ltd., UK) were used to pump the reaction solution into the reactor and remove the effluent. A diagrammatic representation of the experimental setup is shown in Fig. 12. Solution 1 is 0.84M sodium bromate and 0.38M malonic acid. Solution 2 is 1.28M sulphuric acid and 0.22M sodium bromide. The ruthenium catalyst solution was prepared by recrystallizing  $\text{Ru}(\text{bpy})_3\text{SO}_4$  from the chloride salt with sulphuric acid and then dissolving 0.5 g of crystal  $\text{Ru}(\text{bpy})_3\text{SO}_4$  in 18 mL of deionized water to make 0.025M  $\text{Ru}(\text{bpy})_3^{2+}$  solution.

A thin layer chromatography (TLC) plate (T-7270, Sigma-Aldrich, UK) was used as a substrate for the ruthenium catalyst. The silica gel layer was 250  $\mu\text{m}$  thick, the mean particle size was 2–25  $\mu\text{m}$ , and the mean pore diameter was 60  $\text{\AA}$ . As supplied, the plate measured 5  $\times$  20  $\text{cm}^2$  and was divided (by scoring and then snapping over the side of a bench) into sections measuring 5  $\times$  6.5  $\text{cm}^2$  to fit into a Petri dish. A TLC plate section was placed into a Petri dish containing a mixture of 14 mL of deionized water and 0.9 mL of ruthenium catalyst solution. The plate was left in the Petri dish overnight (for approximately 18 h) and then removed and placed into a Petri dish containing deionized water for 30 min to rinse off excess catalyst. The plate was then moved to a third Petri dish containing deionized water for a further 30 min before being used. During the course of a 5-h experimental session the catalyst slowly washed out of the TLC plate and as a consequence the plate became less excitable as the experiment progressed.

The digital projector was a Sanyo PLC-XT21L. The brightness setting of the projector was gradually reduced during the course of a 5-h experimental session to compensate for the washing out of the catalyst, so as to make sure that waves were able to propagate in the fully excitable regime in all dark areas in the projected image, but unable to propagate in light areas.

During the attempts to achieve a uniform subexcitable regime across the whole plate discussed in Sec. II, it was found that the level of excitability varied across the plate. It is believed that these variations were caused by differences in the amount of catalyst in different parts of the plate. However, whereas the range of conditions needed to obtain a subexcitable regime is very narrow, the range needed to obtain a consistent regime in which the substrate is fully excitable in dark areas but unexcitable in light areas is much broader and the behavior of propagating waves in this regime is not very sensitive to small changes in illumination or excitability.

The light intensity in dark areas of the projected image is estimated at less than  $0.1 \text{ mW cm}^{-2}$  and the light intensity in light areas ranged from approximately  $7 \text{ mW cm}^{-2}$  at the beginning of an experimental run to approximately  $4 \text{ mW cm}^{-2}$  at the end of an experimental run. These estimates are based on measurements with a Centronix OSD1-5T photodiode, using the spectral response curve in the manufacturer's datasheet and assuming that the white light produced by the digital projector has a peak wavelength at 550 nm and a power spectrum that is approximately symmetrically distributed around this peak in the range 425–750 nm.

The projected circuit pattern mask was similar to that used during simulation (shown in Fig. 11), but scaled down by a factor of 3. The dimensions of the projected circuit pattern were  $33.5 \times 36 \text{ mm}^2$ . The width of each channel in the projected image was 1.5 mm. Each pixel in the projected image measured  $0.3 \times 0.3 \text{ mm}^2$ .

The period to use for the valves in the projected image (corresponding to  $P$  in the simulation) was determined at the beginning of the experimental session by timing the passage of a wave from one valve to another. A period of 280 s was used, so every valve was open for 140 s, then closed for 140 s, and so on. Images were captured once every 5 s with a Lumenera Infinity 2 USB camera fitted with a 455-nm narrow bandpass interference filter.

## 2. Results

The propagation speed for waves was 6.7 mm/min and it took 10 min for waves to propagate from inputs to outputs. Figure 13 compares snapshots taken from experimental runs with the corresponding snapshots from simulations. The cases  $x = 3, 6,$  and  $15$  are shown. For each case, one pair of images shows the circuit halfway through its operation and another pair of images shows the circuit at the end of its operation. The experimental snapshots were produced by taking the difference of images 30 s apart.

Reference [32] shows seven cases ( $x = 3, 4, 6, 8, 9, 12,$  and  $15$ ) that demonstrate the circuit functioning correctly. Because inputs  $x_0$  and  $x_1$  enter an OR gate immediately, these seven cases fully exercise every gate in the circuit.

The following were the main problems encountered during experiments.

(i) It was not easy to arrange for perfectly synchronized waves to enter the inputs. However, because each valve is open from  $0.25P$  before the time that a wave is expected to pass through it until  $0.25P$  after this time, input waves can pass through the inputs up to 70 s before or after they are expected (corresponding to a distance of  $\pm 7.8 \text{ mm}$ ) before this affects the output of the circuit.

(ii) The wave speed was not identical across the whole plate. In a more complex circuit with many gates cascaded together, small differences in timing will accumulate. Eventually this will lead to some waves being blocked when they should not be or waves passing through a valve when they should be blocked. Assuming that waves travel one unit of length in one unit of time, valves are open at  $t_0 + 2n - 0.5 < t \leq t_0 + 2n + 0.5$  for any integer  $n$  and closed at other times, all input-output paths for a complex circuit are the same length  $L$ , and all input waves enter the circuit at  $t_0$ , then the average speed of a wave must be between  $\frac{L}{L+0.5}$  and  $\frac{L}{L-0.5}$  to avoid any waves being blocked by valves when they should not be and vice versa. No matter how good the experimental conditions are, for a large enough  $L$  this requirement will not be met. It may be possible to overcome this problem by periodically resynchronizing waves to a time reference. One way of doing this is described in Sec. IV F.

## IV. DISCUSSION AND CONCLUSION

We have designed a four-bit input, two-bit output integer square root circuit in BZ excitable media and successfully implemented the circuit experimentally. We believe that one of the reasons why the implementation of this circuit was straightforward is that it makes use of only two of the many phenomena that BZ excitable media can exhibit: constant speed wave propagation and the annihilation of colliding wave fronts. Neither of these phenomena are particularly sensitive to variations in experimental conditions. For the same reason, there is also a reasonably good correspondence between qualitative simulations and experiments.

From a theoretical point of view, it is interesting and useful to consider all possible phenomena in excitable media that can be used for information processing. From a practical point of view, it is easier to implement those systems that depend on the fewest phenomena: Arranging experimental conditions so that all of the phenomena required by a system are present simultaneously is easier if only a few phenomena are required than if many are required. This was borne out in our attempts to use the subexcitable regime for implementing fork structures before switching to the time-dependent wave selection scheme.

Whereas previous work on logic circuits in structured excitable media has used static circuit patterns, we exploit the fact that the circuit pattern is projected using a digital projector to create periodic inhibiting valves, which have the effect of selecting waves based on the time that it takes them to propagate from one valve to another. Because the behavior of a gate in this scheme is dependent only on the distances between terminals rather than on the precise geometry of the gate, paths making up a gate can often be folded to make a compact circuit. An interesting avenue for further research is to explore other ways in which patterns periodically repeating in time can be used in excitable media information processing schemes.

### A. Auxiliary 1 inputs

The most obvious and straightforward way to represent a Boolean value using events is to use the occurrence of an event within a defined period of time to represent a Boolean 1 value and the nonoccurrence of an event within a defined period of time to represent a Boolean 0 value. This is the

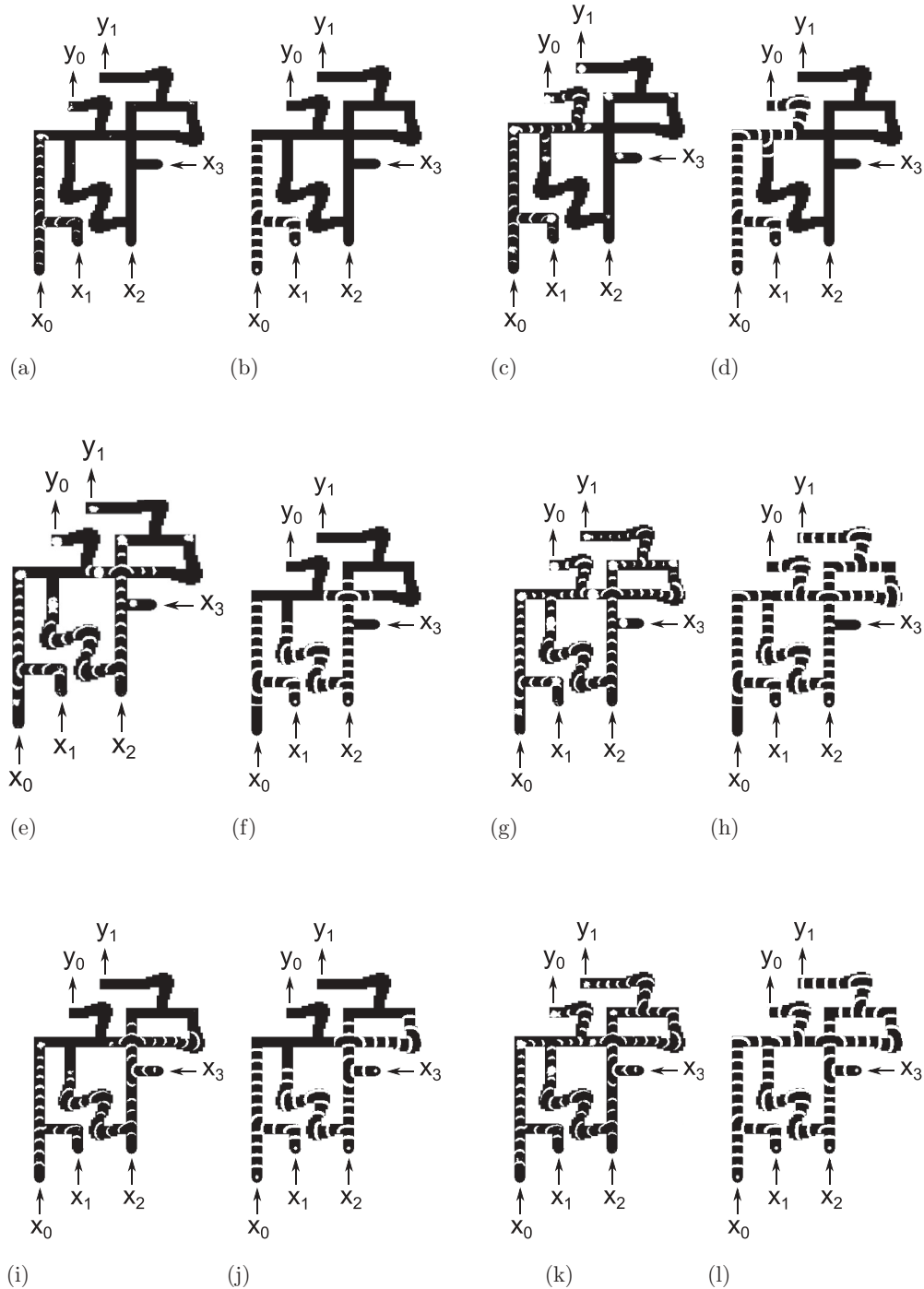


FIG. 13. Experiment and simulation for cases  $x = 3, 6,$  and  $15$  (note that for the  $x = 6$  case no wave emerges from  $y_0$  because the valve at  $y_0$  is closed when a wave reaches it): (a)  $x = 3$  experiment, halfway through; (b)  $x = 3$  simulation, halfway through; (c)  $x = 3$  experiment, complete; (d)  $x = 3$  simulation, complete; (e)  $x = 6$  experiment, halfway through; (f)  $x = 6$  simulation, halfway through; (g)  $x = 6$  experiment, complete; (h)  $x = 6$  simulation, complete; (i)  $x = 15$  experiment, halfway through; (j)  $x = 15$  simulation, halfway through; (k)  $x = 15$  experiment, complete; and (l)  $x = 15$  simulation, complete.

representation that has been used by previous implementations of logic circuits in BZ excitable media and is also used in this paper. One drawback of this method of representing Boolean values is that auxiliary 1 inputs are needed in order to perform certain functions, such as a NOT operation (a NOT gate can be replaced by an AND-NOT gate with a constant auxiliary 1 input fed to the noninverting input). Below we consider other

possible mappings between events and Boolean values that do not have this problem.

### B. Dual-rail logic

One alternative representation is dual-rail logic, in which a Boolean value is represented using a pair of channels. A



Boolean 1 value is represented by the presence of a wave on one channel and a Boolean 0 value is represented by the presence of a wave on the other channel [35]. An obvious disadvantage of a dual-rail representation is that it requires two channels rather than one for each Boolean value. Generally, this disadvantage can be offset against the following potential advantages.

(i) Inversion can be implemented by crossing over two signals. No auxiliary inputs are required.

(ii) Dual-rail logic can facilitate self-timing. A Boolean value is always represented by one event or another: There is no need for an external time reference.

(iii) A set of complete Boolean functions with a dual-rail representation can be built from an incomplete set of single-rail Boolean functions (for example, one that does not have a NOT operation or does not have a complete OR operation) [31].

However, in the context of the work in this paper these do not seem to be significant advantages: The natural AND-NOT behavior of a collision between waves means that inversion is already a straightforward operation (provided that an auxiliary Boolean 1 value is available) and self-timing cannot be fully realized without a mechanism that has a permanent internal state, which is difficult to implement efficiently in BZ excitable media. Since OR and AND-NOT gates (or just AND NOT alone) already form a logically complete set of Boolean functions, the ability to construct a complete set of Boolean functions from an incomplete set is no advantage.

**C. Using phase to represent Boolean values**

Another alternative is to use the phase of a wave to represent Boolean values. By *phase* we mean the delay between a reference time  $t_0$  and the occurrence of a wave in a particular location. Consider a location  $p$  on a channel. The propagation of a wave through  $p$  at some time  $t_0$  (which we will call a leading wave) represents a logic 1 and the propagation of a wave through  $p$  at time  $t_0 + 1$  (which we will call a lagging wave) represents a logic 0.

This representation makes the construction of an OR gate very straightforward. A T shape without any valves can be used as an OR gate because with this representation there is no way that a signal can propagate backward through an input: If both input waves are either leading or lagging then they will merge and the output will follow the inputs, delayed by two time units. Otherwise one input wave will arrive at  $t_0$  and will begin propagating backward through the other input (as well as propagating to the output); however, it will encounter the other input wave, arriving at  $t_0 + 1$ , and will annihilate with it.

A NOT gate in this representation can be constructed as in Fig. 14. For the case when a wave enters the NOT gate at  $a$  at  $t_0$  (when both  $v_1$  and  $v_2$  are open), the wave will reach  $v_1$  at  $t_0 + 1$  when  $v_1$  is closed and will not pass through; however, it will reach  $v_2$  at  $t_0 + 2$  when  $v_2$  is open and will end up exiting the inverter at  $b$  at  $t_0 + 4$ . For the case when a wave enters the NOT gate at  $a$  at  $t_0 + 1$  (when both  $v_1$  and  $v_2$  are closed), the wave will reach  $v_1$  at  $t_0 + 2$  when  $v_1$  is open and will pass through to reach  $b$  at  $t_0 + 3$ . The NOT gate turns a leading wave into a lagging wave and a lagging wave into a leading wave. Between them, the NOT gate and the OR gate can be used to make any desired logic circuit.

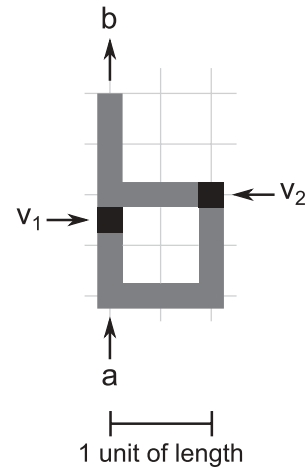


FIG. 14. A NOT gate that turns a leading wave into a lagging wave and a lagging wave into a leading wave.

**D. Crossover**

Although the four-bit integer square root circuit described in this paper does not require any wires to be crossed over, this section describes a structure that may be used to implement crossing wires when this is required. Dewdney [36] showed how to implement a planar crossover circuit using any functionally complete set of logic gates. A crossover circuit can be made from ten AND-NOT gates.

We can do better than this however. Figure 15 shows a crossover that will permit  $x$  to propagate to  $x'$  (but not to  $y$  or  $y'$ ) or alternatively permit  $y$  to propagate to  $y'$  (but not to  $x$  or  $x'$ ) so long as both  $x$  and  $y$  do not occur at the same time. This crossover structure works because the  $x'$  output is the only terminal an even distance from  $x$  and the  $y'$  output is the only terminal an even distance from  $y$ . If we wish to deal with the case when  $x$  and  $y$  may both occur at the same time, we can modify this structure by delaying any wave input at  $y$  by two time units. To keep the output signals synchronized, any wave output at  $x'$  is also delayed by two time units. This is shown in Fig. 15.

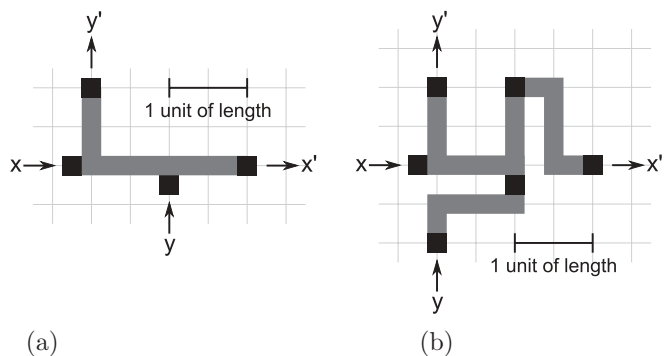


FIG. 15. Structures for crossing over  $x$  and  $y$ : (a) crossover structure for  $x$  and  $y$ , in which  $x$  will propagate to  $x'$  and alternatively  $y$  will propagate to  $y'$ , but the crossover fails if both  $x$  and  $y$  occur at the same time, and (b) crossover structure for  $x$  and  $y$  that will work even if  $x$  and  $y$  occur at the same time.

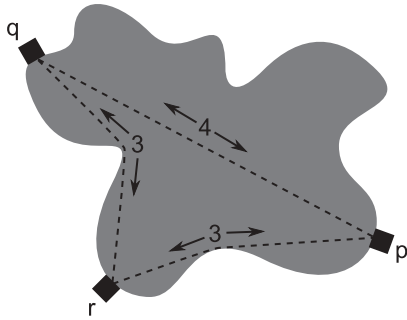


FIG. 16. Arbitrarily shaped region of excitable media with three terminals with valves. Distances between pairs of valves are shown.

### E. General abstract description of time-dependent wave selection

The reason why we are able to make such a simple crossover mechanism (which is much simpler than one made from ten AND-NOT gates) is that we are working at a level of abstraction lower than that of Boolean logic gates. Time-dependent wave selection can be used to implement logic gates, but it can also be used to implement other structures more directly than via logic gates.

A general abstract description of the behavior of waves propagating between valves in a region of BZ excitable medium is stated below. The internal distance  $|\vec{p}q|$  between any two points  $p$  and  $q$  within a region is the length of the shortest path between  $p$  and  $q$  that lies within the region. Here we assume that the internal distance between any pair of terminals is an integral length and that the region contains no unexcitable holes.

A region may be of any shape and may contain any number of terminals on its boundary, with each terminal having a valve that is open only at even times. A wave starting from a given terminal  $p$  at an even time  $t$  will propagate to a terminal  $q$  at time  $t + |\vec{p}q|$  so long as the distance  $|\vec{p}q|$  is even and there is no wave that begins from any terminal  $r$  at an even time between  $t - |\vec{r}p|$  and  $t + |\vec{p}q| - |\vec{r}q|$  (where this interval includes the former time but not the latter) and that propagates uninterrupted to  $p$ . This is illustrated in Fig. 16. In Fig. 16 a wave from  $p$  at  $t_0$  will propagate to  $q$  at  $t_0 + 4$  unless there is a wave from  $r$  at  $t_0 - 2$  or  $t_0$ .

To simplify the process of designing structures, in this paper we have used only regions with edge lengths that are multiples of half a unit length and any pair of edges is either parallel or orthogonal. To reduce the space that structures occupy, we have used thin regions.

### F. Resynchronization

In Sec. III B2 we discussed the need to resynchronize waves to a time reference to remove small variations in wave propagation time that may accumulate for large logic circuits. Figure 17 shows one possible mechanism for doing this. This mechanism uses the time of valve opening as a

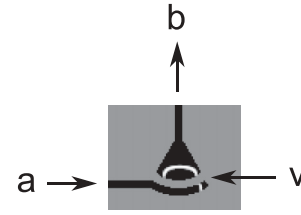


FIG. 17. Mechanism that will synchronize waves with valve timing. All points along valve  $v$  are the same distance from  $b$  so that a wave originating from  $a$  will reach  $b$  at a constant time after  $v$  opens, regardless of the exact time at which it entered  $a$ .

reference time. Waves exiting the mechanism are always at a known offset from the valve opening time. The entire curved region labeled  $v$ , between the two dark regions, is a valve in the shape of a circular arc. When valve  $v$  is closed, a wave entering  $a$  will propagate along the curved channel beneath  $v$ . When the valve opens, the wave will propagate through the valve and head toward  $b$ . Because all points along valve  $v$  are equidistant from  $b$ , the wave will reach  $b$  at a constant time after  $v$  opens, regardless of the exact time at which it entered  $a$ .

### G. Parallels with neural information processing

The phenomena used to implement the scheme described here have some similarities to phenomena that are present in neural circuits. Both involve the propagation of excitations at a predictable speed and have regions in which excitations can be inhibited. Among the large number of things that we do not yet understand about neural information processing is the range of roles that inhibitory synapses play. The scheme described here suggests an alternative possibility: It may be possible to use neurons to construct circuits in which thresholding plays no direct role in information processing. In such a scheme neurons serve simply as repeaters and as the providers of a medium in which excitations can propagate and interact. Periodically spiking inhibitory neurons with axons leading to inhibitory synapses on other neurons could play the same role that valves play in this paper.

During the evolution of nervous systems, we do not know whether the physiology of neurons and the phenomena that are exhibited by neurons were selected because they are well suited to their role as components of evolvable, adaptable information processing systems or whether their evolution was highly constrained by other factors. If we assume the former, we should consider systematically evaluating models that incorporate one or more of the phenomena exhibited by neurons for their potential use in evolutionary computing schemes (as has been done in Refs. [37,38]). The work in this paper is one such model.

### ACKNOWLEDGMENTS

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